# First passage time problems in stochastic hybrid systems

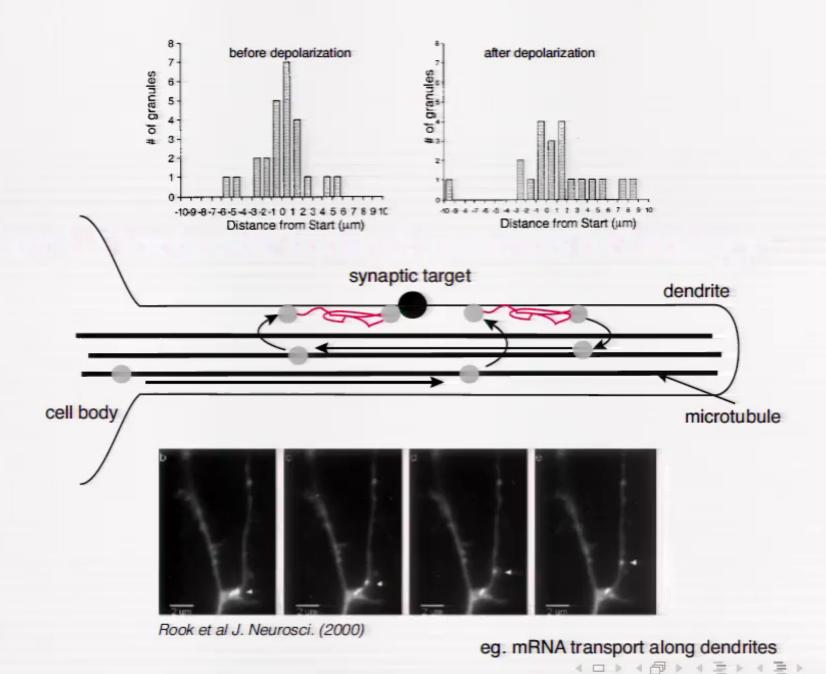
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# [A] BIDIRECTIONAL MOTOR-DRIVEN INTRACELLULAR TRANSPORT

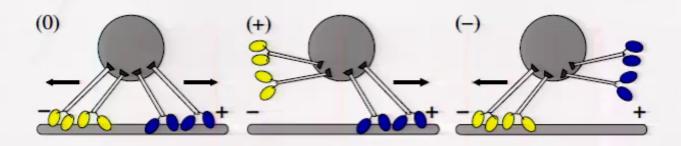


# [A] TUG-OF-WAR MODEL

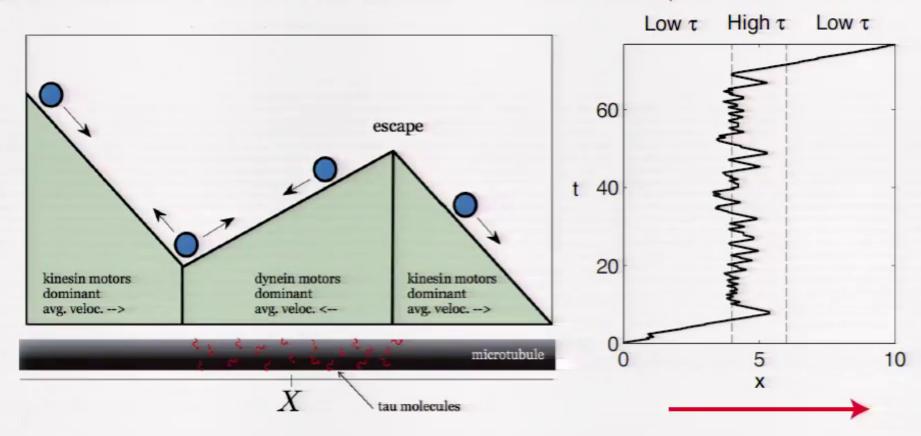
- Let X(t) denote position of motor-cargo complex on MT track.
- Discrete state given by  $(n_+, n_-)$  where  $n_+$   $(n_-)$  is number of kinesin (dynein) motors bound to MT
- Let  $v(n_+, n_-)$  be velocity of motor complex in state  $(n_+, n_-)$
- Piecewise deterministic dynamics of motor complex is

$$\frac{dX}{dt} = v(n_+, n_-)$$

 Transitions between different states described by a jump Markov process (Klumpp et al)



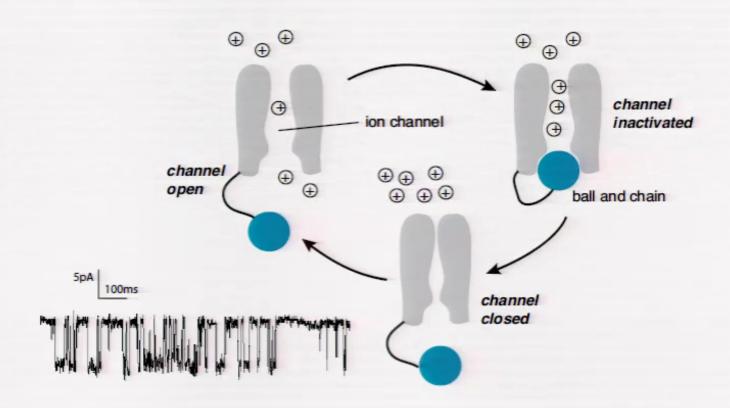
## [A] LOCAL SIGNALING AND AN ESCAPE PROBLEM (NEWBY/PCB 2010)



- MAPs (tau, MAP2) can bind to microtubules and reduce the binding rate of kinesin
- Provides a possible explanation for experimentally observed oscillatory motion of motor complexes
- Now have X-dependent velocities  $v = v(n_+, n_-, X)$  and transition rates.

## [B] STOCHASTIC ION CHANNELS AND VOLTAGE FLUCTUATIONS

 Single ion channels fluctuate rapidly between open and closed states in a stochastic fashion.



# [B] STOCHASTIC CONDUCTANCE-BASED MODEL (PCB/Newby/Keener 2013)

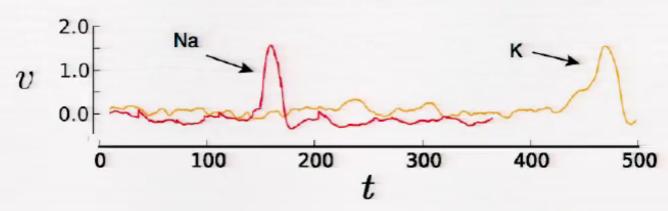
- Suppose a neuron has  $n \le N$  open Na<sup>+</sup> channels and  $m \le M$  open K<sup>+</sup> channels
- $\bullet$  Voltage V(t) evolves according to piecewise deterministic dynamics

$$\frac{dv}{dt} = F(v, m, n) \equiv \frac{n}{N} f_{Na}(v) + \frac{m}{M} f_K(v) - g(v).$$
with  $f_i(v) = \bar{g}_i(v_i - v)$ 

Assume each channel satisfies the simple kinetic scheme

$$C(closed) \xrightarrow{\alpha_i(v)} O(open), \quad i = \text{Na, K},$$
  
 $\beta_i(v)$ 

Ion channel fluctuations can induce spontaneous action potentials.



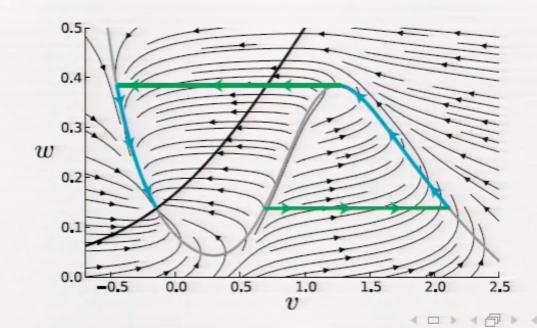
#### [B] MORRIS-LECAR MODEL OF NEURAL EXCITABILITY

• In the limit of fast Na<sup>+</sup> channels and infinite K<sup>+</sup> channels ( $M \to \infty$ ) we obtain the deterministic Morris-Lecar (ML) model

$$\frac{dv}{dt} = \frac{\alpha_{Na}(v)}{\alpha_{Na}(v) + \beta_{Na}(v)} f_{Na}(v) + w f_{K}(v) - g(v)$$

$$\frac{dw}{dt} = \alpha_{K}(v)(1 - w) - \beta_{K}(v)w,$$

- Examine excitability using slow/fast analysis
- Require large perturbations (rare events) to induce an action potential

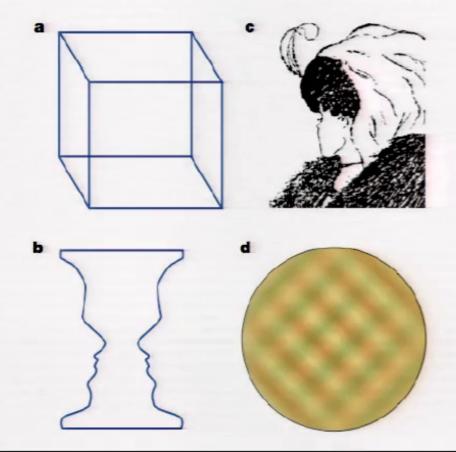


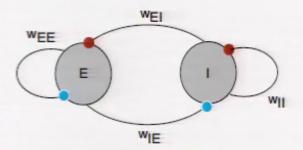
## [C] GENETIC SWITCHES/BISTABLE NEURAL NETWORKS

continuous variable = promoter protein concentration

discrete variable = state of promoter



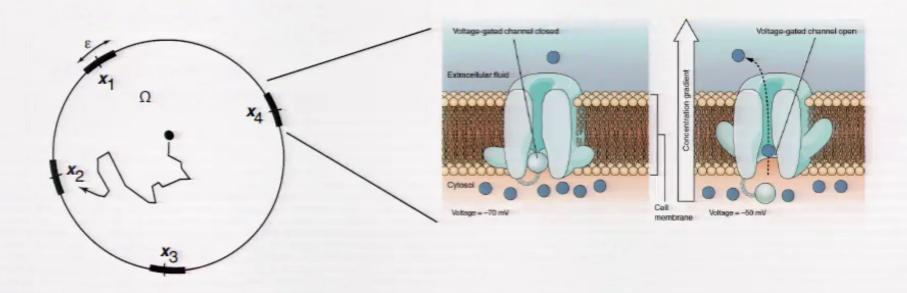




continuous variables = synaptic currents in each population

discrete variables= number of spiking neuronsin each population

# [D] ESCAPE FROM SUBCELLULAR DOMAINS WITH SWITCHING BOUNDARIES



- Diffusion of molecules within a subcellular domain  $\Omega$  with stochastic channels in the membrane  $\partial\Omega$
- Consider N narrow gates  $\partial \Omega_k^{\varepsilon}$ ,  $k \in \{1, ..., N\}$ , with the kth gate given by the  $\varepsilon > 0$  neighborhood of  $\mathbf{x}_k \in \partial \Omega$  defined according to

$$\partial \Omega_k^{\varepsilon} := \{ \mathbf{x} \in \partial \Omega : |\mathbf{x} - \mathbf{x}_k| < \varepsilon \}.$$

• Let  $\mathbf{n}(t) \in \{0,1\}^N$  be an irreducible Markov process whose k-th component,  $n_k(t) \in \{0,1\}$ , controls the state of the k-th gate.

Part II. Analysis of first passage time problems

#### 1D STOCHASTIC HYBRID SYSTEM

Consider the 1D system

$$\frac{dx}{dt} = \frac{1}{\tau_x} F_n(x), \quad x \in \mathbb{R}, \quad n = 1, \dots, K$$

- Jump Markov process  $m \to n$  with transition rates  $W_{nm}(x)/\tau_n$ .
- Set  $\tau_x = 1$  and introduce the small parameter  $\epsilon = \tau_n/\tau_x$
- Chapman-Kolmogorov (CK) equation for  $p_n(x,t) = \mathbb{E}[p(x,t)1_{n(t)=n}]$  is

$$\frac{\partial p_n}{\partial t} = -\frac{\partial [F_n(x)p_n(x,t)]}{\partial x} + \frac{1}{\epsilon} \sum_{m=1}^K A_{nm}(x)p_m(x,t)$$

where

$$A_{nm}(x) = W_{nm}(x) - \sum_{k=1}^{K} W_{kn}(x) \delta_{m,n}.$$

• Assume that there exists a unique stationary density  $\rho_n(x)$  with

$$\sum_{m} A_{nm}(x) \rho_m(x) = 0$$

#### QUASI-STEADY-STATE APPROXIMATION (PCB/NEWBY)

• In the limit  $\epsilon \to 0$ , obtain mean-field equation

$$\frac{dx}{dt} = \mathcal{F}(x) \equiv \sum_{n=1}^{K} F_n(x) \rho_n(x),$$

Decompose the probability density as

$$p(x,n,t) = C(x,t)\rho_n(x) + \epsilon w_n(x,t),$$
 where  $\sum_n p_n(x,t) = C(x,t)$  and  $\sum_n w_n(x,t) = 0$ .

• Asymptotic expansion in  $\epsilon$  yields FP equation

$$\frac{\partial C}{\partial t} = -\frac{\partial}{\partial x}(\mathcal{F}C) + \epsilon \frac{\partial}{\partial x} \left( \mathcal{D} \frac{\partial C}{\partial x} \right),$$

Drift term given by mean-field equation, and diffusion coefficient

$$\mathcal{D}(x) = \sum_{n} Z_n(x) F_n(x),$$

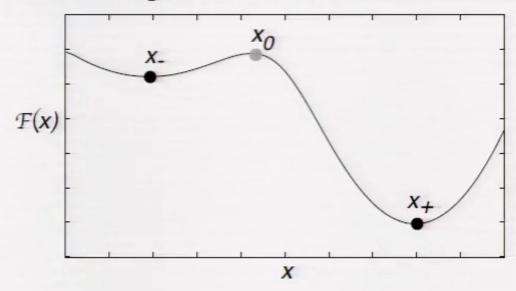
where

$$\sum_{m} A_{nm}(x) Z_m(x) = -[\mathcal{F}(x) - F_n(x)] \rho_n(x)$$

7

## FIRST-PASSAGE TIME (FTP) PROBLEM I

Suppose that mean field equation is bistable



- Let T(x) be the stochastic time for the particle to exit at  $x_0$  starting at x
- Introduce the survival probability  $\mathbb{P}(x,t)$  that the particle has not yet exited at time t:

$$\mathbb{P}(x,t) = \int_0^{x_0} \sum_n p_n(x',t|x,0) dx'.$$

and define the first passage time (FPT) density

$$f(x,t) = -\frac{\partial \mathbb{P}(x,t)}{\partial t}.$$

# FIRST-PASSAGE TIME (FTP) PROBLEM II

• The mean first passage time (MFPT)  $\tau(x)$  is

$$\tau(x) = \langle T(x) \rangle \equiv \int_0^\infty f(x,t)tdt = \int_0^\infty \mathbb{P}(x,t)dt,$$

• In limit  $\epsilon \to 0$ , expect MFPT to have the Arrhenius-like form

$$\tau(x_{-}) = \frac{2\pi\Gamma(x_{0}, x_{-})}{\sqrt{|\Phi''(x_{0})|\Phi''(x_{-})}} e^{[\Phi(x_{0}) - \Phi(x_{-})]/\epsilon}.$$

where  $\Phi(x)$  is a **quasipotential** and  $\Gamma$  is a prefactor.

QSS approximation yields the approximate quasipotential

$$\Phi_{\rm QSS}(x) \equiv -\int^x \frac{\mathcal{F}(x')}{\mathcal{D}(x')} dx'$$

- ⇒ may generate exponentially large errors in MFPT
- FP equation is 2nd-order, whereas CK equation is K-th order ⇒ QSS can breakdown at boundaries

#### PATH-INTEGRAL REPRESENTATION (PCB/Newby 2014)

- Determine  $\Phi(x)$  using large deviation theory/path integrals/WKB
- Consider the eigenvalue equation

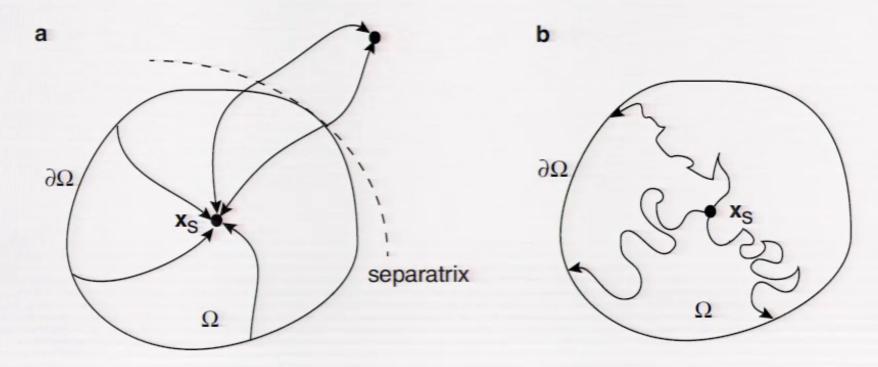
$$\sum_{m} [A_{nm}(x) + q\delta_{n,m}F_{m}(x)] R_{m}^{(s)}(x,q) = \lambda_{s}(x,q)R_{n}^{(s)}(x,q),$$

and let  $\xi_m^{(s)}$  be the adjoint eigenvector.

- Perron-Frobenius theorem shows that there exists a real, simple Perron eigenvalue labeled by s = 0, say, such that  $\lambda_0 > \text{Re}(\lambda_s)$  for all s > 0
- Path-integral representation of PDF

$$P(x,\tau) = \int_{x(0)=x_*}^{x(\tau)=x} \exp\left(-\frac{1}{\epsilon} \int_0^{\tau} [p\dot{x} - \lambda_0(x,p)]dt\right) \mathcal{D}[p]\mathcal{D}[x]$$

#### "ZERO ENERGY" PATHS



- (a) Deterministic trajectories converging to a stable fixed point  $x_s$ . Boundary of basin of attraction formed by a union of separatrices
- (b) Noise-induced paths of escape

#### MEAN-FIELD EQUATIONS

• We have the trivial solution p = 0 and  $R_m^{(0)}(x, 0) = \rho_m(x)$  with

$$\sum_{m} A_{nm}(x) \rho_m(x) = 0$$

• Differentiating the eigenvalue equation with respect to p and then setting p = 0,  $\lambda_0 = 0$  shows that

$$\frac{\partial \lambda_0(x,p)}{\partial p}\bigg|_{p=0} \rho_n(x) = F_n(x)\rho_n(x) + \sum_m A_{nm}(x) \left. \frac{\partial R_m^{(0)}(x,p)}{\partial p} \right|_{p=0}$$

• Summing both sides wrt *n* and using  $\sum_{n} A_{nm} = 0$ ,

$$\frac{\partial \lambda_0(x)}{\partial p}\bigg|_{p=0} = \sum_n F_n(x) \rho_n(x)$$

• Hamilton's equation  $\dot{x} = \partial \lambda_0(x, p)/\partial p$  recovers mean-field equation

$$\dot{x} = \sum_{n} F_n(x) \rho_n(x).$$

#### MAXIMUM-LIKELIHOOD PATHS OF ESCAPE

• Unique non-trivial solution  $p = \mu(x)$  with positive eigenvector  $R_m^{(0)}(x, \mu(x)) = \psi_m(x)$ :

$$\sum_{m} \left[ A_{nm}(x) + \mu(x) \delta_{n,m} F_m(x) \right] \psi_m(x) = 0$$

• Yields quasipotential  $\Phi(x)$  with  $\Phi'(x) = \mu(x)$  and

$$S[x,p] \equiv \int_{-\infty}^{\tau} \left[ p\dot{x} - \lambda_0(x,p) \right] dt = \int_{x_s}^{x} \Phi'(x) dx.$$

 Equivalent to WKB quasipotential obtained using ansatz for quasistationary solutions

$$p_n(x) = R_n(x) \exp\left(-\frac{1}{\epsilon}\Phi(x)\right),$$

Part III. Stochastic ion-channels revisited

#### REDUCED MORRIS-LECAR MODEL

• Let n, n = 0, ..., N be the number of open sodium channels:

$$\frac{dv}{dt} = F_n(v) \equiv \frac{1}{N} f(v) n - g(v),$$
 with  $f(v) = g_{Na}(V_{Na} - v)$  and  $g(v) = -g_{eff}[V_{eff} - v] + I_{ext}$ .

 The opening and closing of the ion channels is described by a birth-death process according to

$$n \rightarrow n \pm 1$$
,

with rates

$$\omega_{+}(n) = \alpha(v)(N-n), \quad \omega_{-}(n) = \beta n$$

Take

$$\alpha(v) = \beta \exp\left(\frac{2(v - v_1)}{v_2}\right)$$

for constants  $\beta$ ,  $v_1$ ,  $v_2$ .

## CHAPMAN-KOLMOGOROV EQUATION

CK equation is

$$\frac{\partial p_n}{\partial t} = -\frac{\partial [F_n(v)p_n(v,t)]}{\partial v} + \frac{1}{\epsilon} \sum_{n'} A_{nm}(v)p_m(v,t),$$

$$A_{n,n-1} = \omega_+(n-1), A_{nn} = -\omega_+(n) - \omega_-(n), A_{n,n+1} = \omega_-(n+1).$$

• There exists a unique steady state density  $\rho_n(v)$  for which

$$\sum_{m} A_{nm}(v) \rho_m(v) = 0$$

where

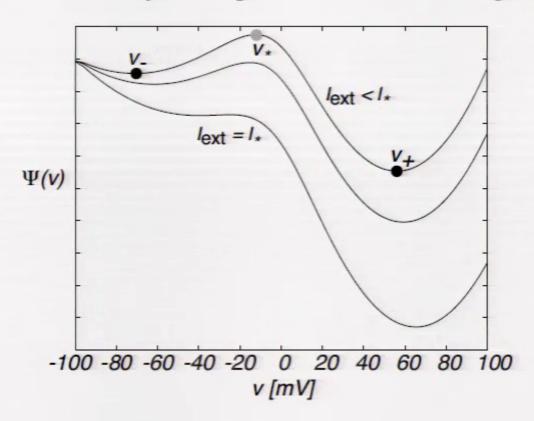
$$\rho_n(v) = \frac{N!}{(N-n)!n!} a(v)^n b(v)^{N-n}, \quad a(v) = \frac{\alpha(v)}{\alpha(v) + \beta}, \ b(v) = 1 - a(v).$$

#### MEAN-FIELD LIMIT

• In the limit  $\epsilon \to 0$ , we obtain the mean-field equation

$$\frac{dv}{dt} = \sum_{n} F_n(v) \rho_n(v) = a(v) f(v) - g(v) \equiv -\frac{d\Psi}{dv},$$

Assume deterministic system operates in a bistable regime



#### PERRON EIGENVALUE I

• Eigenvalue equation for  $\lambda_0$  and  $R^{(0)} = \psi$ :

$$(N - n + 1)\alpha\psi_{n-1} - [\lambda_0 + n\beta + (N - n)\alpha]\psi_n + (n + 1)\beta\psi_{n+1}$$
$$= -p\left(\frac{n}{N}f - g\right)\psi_n$$

Consider the trial solution

$$\psi_n(x,p) = \frac{\Lambda(x,p)^n}{(N-n)!n!},$$

• Yields the following equation relating  $\Lambda$  and  $\mu$ :

$$\frac{n\alpha}{\Lambda} + \Lambda\beta(N-n) - \lambda_0 - n\beta - (N-n)\alpha = -p\left(\frac{n}{N}f - g\right).$$

Collecting terms independent of n and terms linear in n yields

$$p = -\frac{N}{f(x)} \left( \frac{1}{\Lambda(x,p)} + 1 \right) (\alpha(x) - \beta(x)\Lambda(x,p)),$$

and

$$\lambda_0(x,p) = -N(\alpha(x) - \Lambda(x,p)\beta(x)) - pg(x).$$

## THE QUASIPOTENTIAL

Non-trivial solution ryields

$$p = \mu(x) \equiv N \frac{\alpha(x)f(x) - (\alpha(x) + \beta)g(x)}{g(x)(f(x) - g(x))}.$$

ullet The corresponding quasipotential  $\Phi$  is given by

$$\Phi(x) = \int_{-\infty}^{x} \mu(y) dy.$$

Analogous result in full ML model

#### STOCHASTIC MORRIS-LECAR MODEL

• Take  $n \le N$  open Na<sup>+</sup> channels and  $m \le M$  open K<sup>+</sup> channels:

$$\frac{dv}{dt} = F(v, m, n) \equiv \frac{n}{N} f_{Na}(v) + \frac{m}{M} f_K(v) - g(v).$$

Each channel satisfies the kinetic scheme

$$C \underset{\beta_i(v)}{\overset{\alpha_i(v)}{\rightleftharpoons}} O, \quad i = \text{Na, K},$$

- The Na<sup>+</sup> channels fast relative to voltage and K<sup>+</sup> dynamics.
- Chapman–Kolmogorov (CK) equation,

$$\frac{\partial p}{\partial t} = -\frac{\partial (Fp)}{\partial v} + \mathbb{L}_{K}p + \mathbb{L}_{Na}p.$$

• The jump operators  $\mathbb{L}_j$ , j = Na, K, are defined according to

$$\mathbb{L}_j = (\mathbb{E}_n^+ - 1)\omega_j^+(n) + (\mathbb{E}_n^- - 1)\omega_j^-(n),$$
 with  $\mathbb{E}_n^{\pm} f(n) = f(n \pm 1)$ ,  $\omega_i^-(n) = n\beta_i$  and  $\omega_i^+(n) = (N - n)\alpha_i(v)$ .

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#### WKB APPROXIMATION

Introduce quasistationary solution of the form

$$\varphi(v, w, n) = R_n(v, w) \exp\left(-\frac{1}{\epsilon}\Phi(v, w)\right),$$

where  $\Phi(v, w)$  is the **quasipotential** 

To leading order,

$$[\mathbb{L}_{Na} + p_v + h(v, w, p_w)] R_n(v, w) = 0,$$

where

$$p_v = \frac{\partial \Phi}{\partial v}, \quad p_w = \frac{\partial \Phi}{\partial w}$$

and

$$h(v, w, p_w) = \frac{\beta_K}{M\lambda_M} \left[ (e^{-\lambda_M p_w} - 1)\omega_K^+(Mw, v) + (e^{\lambda_M p_w} - 1)\omega_K^-(Mw, v) \right]$$