Optimizing the Kelvin force in a moving target subdomain

Harbir Antil

Department of Mathematical Sciences George Mason University hantil@gmu.edu

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Collaborators: R.H. Nochetto (UMD) and P. Venegas (Chile)

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Motivation

The magnetic field exerts a force on magnetic materials such as magnetic nanoparticles (MNPs). MNPs under the action of external magnetic field are used in:

- medical sciences:
 - as contrast agents to enhance the contrast in MRI
 - as carriers for targeted drug delivery, for instance, to treat cancer cells, tumors (< 0.1% is taken by tumor cells)
 - in gene therapy
 - in magnetized stem-cells
- magnetic tweezers
- Iab-on-a-chip systems that include magnetic particles or fluids
- magnetofection a transfection method



Magnetic drug targeting (MDT)





Experimental setup ([Shapiro et al., 2013])



What are engineers interested in?



> Pull in or attract particles (left): there are already human trials on this.

- Key difficulty: To push or to control particles (center and right). In region A the force is pointing outward, allowing us to push particles.
- The success of the aforementioned applications highly depend on the accurate control of the magnetic force.
- ► Goal: how to approximate a desired magnetic force f by a fixed configuration of magnetic field sources.
- Approach:

$$\min_{\mathbf{F}} \int_0^T \|\mathbf{F} - \mathbf{f}\|_{\mathrm{L}^2(D)}^2 \, dt \quad \text{for } T > 0 \text{ and } D \subset \mathbb{R}^d, d = 2, 3.$$



How does magnetic field manipulate MNPs?

Magnetic force: Magnetic field gradient is required to exert a force and such a force is given by [Rosensweig '97]:

$$\mathbf{F} = (\mathbf{m} \cdot \nabla) \mathbf{H}.$$

► A simplification: After some simplifications (weakly diamagnetic medium, no current sources) and using curl H = 0, the force on a single MNP

$$\mathbf{F} = \frac{V_m \Delta \chi}{2} \nabla |\mathbf{H}|^2$$

 V_m : is the volume of the particle $\Delta \chi = \chi_p - \chi_m$: effective susceptibility.

Fundamental difficulty: magnetic field intensity H is not parallel to the magnetic force F.



Minimization problem and control action

- $\Omega \subset \mathbb{R}^d$, d = 2, 3 open, bounded $D_t \subset \Omega$: time dependent subdomain.
- \blacktriangleright Maxwell's equations: We consider magnetic sources outside Ω then

$$\operatorname{curl} \mathbf{H} = \mathbf{0}, \quad \operatorname{div} \mathbf{H} = 0 \quad \text{in } \Omega.$$

Dipole approximation:

$$\mathbf{H}(\boldsymbol{x},t) = \sum_{i=1}^{n_p} \alpha_i(t) \left(d \frac{(\boldsymbol{x} - \boldsymbol{x}_i)(\boldsymbol{x} - \boldsymbol{x}_i)^\top}{|\boldsymbol{x} - \boldsymbol{x}_i|^2} - \mathbb{I} \right) \frac{\widehat{\mathbf{d}}_i}{|\boldsymbol{x} - \boldsymbol{x}_i|^d} = \sum_{i=1}^{n_p} \alpha_i(t) \mathbf{H}_i(\boldsymbol{x})$$

 $oldsymbol{x}_i \in \mathbb{R}^d \setminus \overline{\Omega}$: dipole positions $\widehat{\mathbf{d}}_i \in \mathbb{R}^d$: field direction







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Problem 1: Fixed final time

Minimization problem:

$$\begin{split} & \min_{\boldsymbol{\alpha}\in\mathcal{H}_{ad}}\mathcal{J}(\boldsymbol{\alpha}), \quad \text{with } \frac{1}{2}\int_{0}^{T}\|\nabla|\mathbf{H}(\boldsymbol{\alpha})|^{2} - \mathbf{f}\|_{\mathrm{L}^{2}(D_{t})}^{2}dt + \frac{\lambda}{2}\int_{0}^{T}|d_{t}\boldsymbol{\alpha}|^{2}dt, \\ & \text{with } \boldsymbol{\alpha}(t) \coloneqq (\alpha_{1}(t), \dots, \alpha_{n_{p}}(t))^{\top} \text{ in} \\ & \mathcal{H}_{ad} \coloneqq \left\{\boldsymbol{\alpha}\in[\mathrm{H}^{1}(0,T)]^{n_{p}}:\boldsymbol{\alpha}(0) = \boldsymbol{\alpha}_{0} \quad \text{and} \quad \boldsymbol{\alpha}_{*}\leq\boldsymbol{\alpha}(t)\leq\boldsymbol{\alpha}^{*}, \; \forall t\in[0,T]\right\}. \end{split}$$

Reformulation:

$$\min_{\boldsymbol{\alpha}\in\mathcal{H}_{ad}}\mathcal{J}(\boldsymbol{\alpha}), \quad \text{with } \mathcal{J}(\boldsymbol{\alpha}) = \frac{1}{2}\int_0^T \left(\sum_{i=1}^d \|\boldsymbol{\alpha}^\top \mathbf{P}_i \boldsymbol{\alpha} - \mathbf{f}_i\|_{\mathrm{L}^2(D_t)}^2 + \lambda |d_t \boldsymbol{\alpha}|^2\right) dt.$$



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Problem 2: Minimizing the final time

- ▶ Unknown f. Since the final time *T_F* is an unknown, thus f is not meaningful quantity. We treat f as an unknown.
- Let $\rho \in C^1[0, s_F]$ be a parameterization of C with respect to arc length.



• Assume that the barycenter $x_C(t)$ of D_t moves along curve C at speed $\theta(t) > 0$ with initial position x_I and final x_F such that

$$\int_0^{T_F} \theta(\tau) d\tau = s_F.$$

 \blacktriangleright We define the map $\sigma(\cdot):[0,T_F]\rightarrow [0,s_F]$ as

$$s = \sigma(t) = \int_0^t \theta(\tau) d\tau.$$

• Whence
$$\boldsymbol{x}_{C}(\cdot) = \boldsymbol{\rho} \circ \sigma(\cdot)$$
. Also $d_t \boldsymbol{x}_{C}(t) = \theta(t) d_t \boldsymbol{\rho}(\sigma(t))$.



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Problem 2: Minimizing the final time

$$\min_{(\boldsymbol{\alpha},\theta)\in\mathcal{H}_{ad}\times\mathcal{V}_{ad}}\mathcal{F}(\boldsymbol{\alpha},\theta)$$

with

$$\begin{split} \mathcal{F}(\boldsymbol{\alpha}, \boldsymbol{\theta}) &\coloneqq \int_{0}^{s_{F}} \left(\frac{1}{2\theta(s)} \sum_{i=1}^{d} \|\boldsymbol{\alpha}(s)^{\top} \mathbf{P}_{i} \boldsymbol{\alpha}(s) - \boldsymbol{\rho}_{i}'(s) \boldsymbol{\theta}(s) \|_{\mathrm{L}^{2}(D_{s})}^{2} \\ &+ \frac{\beta}{\theta(s)} + \frac{\lambda}{2} |d_{s} \boldsymbol{\alpha}(s)|^{2} + \frac{\eta}{2} |d_{s} \boldsymbol{\theta}(s)|^{2} \right) ds \end{split}$$

and

$$\begin{split} \mathcal{U}_{ad} &:= \left\{ \boldsymbol{\alpha} \in \mathrm{H}^{1}(0, s_{F}) : \boldsymbol{\alpha}(0) = \boldsymbol{\alpha}_{0} \quad \text{and} \quad \boldsymbol{\alpha}_{*} \leq \boldsymbol{\alpha}(s) \leq \boldsymbol{\alpha}^{*}, \ \forall s \in [0, s_{F}] \right\}, \\ \mathcal{V}_{ad} &:= \left\{ \boldsymbol{\theta} \in \mathrm{H}^{1}(0, s_{F}) : \boldsymbol{\theta}(0) = \boldsymbol{\theta}_{0} \quad \text{and} \quad 0 < \boldsymbol{\theta}_{*} \leq \boldsymbol{\theta}(s) \leq \boldsymbol{\theta}^{*}, \ \forall s \in [0, s_{F}] \right\}. \end{split}$$



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Well-posedness of continuous problem

- ▶ Existence of solution. If $\mathbf{f} \in [L^2(0,T; L^2(\Omega))]^d$ then using direct method of calculus of variations, there exist a solution to Problem 1. Same holds for Problem 2 if $\boldsymbol{\rho}$ is C^1 .
- First order necessary optimality conditions.
 - If $ar{oldsymbol{lpha}}\in\mathcal{U}_{ad}$ solves Problem 1 then

$$\mathcal{J}'(\bar{\boldsymbol{\alpha}})(\boldsymbol{\alpha}-\bar{\boldsymbol{\alpha}})\geq 0 \qquad \forall \boldsymbol{\alpha}\in \mathcal{H}_{ad}.$$

▶ If $(\bar{\alpha}, \bar{\theta}) \in \mathcal{U}_{ad} \times \mathcal{V}_{ad}$ solves Problem 2 then

$$\nabla \mathcal{F}(\bar{\boldsymbol{\alpha}}, \bar{\theta})(\delta \boldsymbol{\alpha}, \delta \theta) \ge 0 \qquad \forall (\boldsymbol{\alpha}, \theta) \in \mathcal{U}_{ad} \times \mathcal{V}_{ad}$$

where $\delta \alpha = \alpha - \bar{\alpha}$, $\delta \theta = \theta - \bar{\theta}$.

Second order sufficient condition. Under the assumption

$$\mathcal{J}''(\bar{\boldsymbol{\alpha}})(\delta \boldsymbol{\alpha})^2 \geq \omega |\delta \boldsymbol{\alpha}|^2_{\mathrm{H}^1(0,T)} \qquad \forall \delta \boldsymbol{\alpha} \in \mathcal{A}(\bar{\boldsymbol{\alpha}})$$

where $\mathcal{A}(\boldsymbol{\alpha}) := \left\{ \mathbf{h} \in \mathrm{H}_{0}^{1}(0,T) : \boldsymbol{\alpha} + \zeta \mathbf{h} \in \mathcal{H}_{ad}, \zeta > 0 \right\}$ there exists a local unique solution to Problem 1. Same applies to Problem 2.



Problem 1: Reference domain

▶ Reference domain. We define a reference domain $\widehat{D} \subset \mathbb{R}^d$ and a map $\mathbf{X} : [0,T] \times \overline{\widehat{D}} \to \overline{\Omega}$, such that for all $t \in [0,T]$

$$\begin{split} \mathbf{X}(t,\cdot) : & \overline{\widehat{D}} \to \overline{D}_t \\ & \widehat{\boldsymbol{x}} \to \mathbf{X}(t, \widehat{\boldsymbol{x}}) = \boldsymbol{\psi}(t) + \boldsymbol{\psi}(t) \widehat{\boldsymbol{x}}, \end{split}$$

 $\pmb{\psi}:[0,T]\to\mathbb{R}^d,\,\psi:[0,T]\to(0,+\infty),$ and $\pmb{\psi},\psi\in\mathrm{H}^1(0,T).$ Moreover, $\widehat{D}=D_0.$

• Reference domain cost. Then, we rewrite \mathcal{J} as

$$\mathcal{J}(\boldsymbol{\alpha}) = \frac{1}{2} \sum_{i=1}^{d} \int_{0}^{T} \|\boldsymbol{\alpha}^{\top} \widehat{\mathbf{P}}_{i} \boldsymbol{\alpha} - \widehat{\mathbf{f}}_{i} \|_{\mathbf{L}^{2}(\widehat{D})}^{2} + \frac{\lambda}{2} \int_{0}^{T} |d_{t} \boldsymbol{\alpha}|^{2} = \mathcal{J}^{1}(\boldsymbol{\alpha}) + \mathcal{J}^{2}(\boldsymbol{\alpha})$$

with $\widehat{\mathbf{P}}_i(t, \widehat{\boldsymbol{x}}) := \mathbf{P}_i(\mathbf{X}(t, \widehat{\boldsymbol{x}}))\psi(t)^{d/2}$ and $\widehat{\mathbf{f}}_i(t, \widehat{\boldsymbol{x}}) = v_i(t, \mathbf{X}(t, \widehat{\boldsymbol{x}}))\psi(t)^{d/2}$, for $i = 1, \dots, d$.



Problem 1: Time discretization

• Let us fix $N \in \mathbb{N}$ and let $\tau := T/N$ be the time step. Now, for $n = 1, \ldots, N$, we define $t^n := n\tau$, $\widehat{\mathbf{P}}_i^n = \widehat{\mathbf{P}}_i(t^n)$ and $\widehat{\mathbf{f}}_i^n$ to be

$$\widehat{\mathbf{f}}_{i}^{n}(\cdot) = \frac{1}{\tau} \int_{t^{n-1}}^{t^{n}} \widehat{\mathbf{f}}_{i}(t, \cdot) dt, \qquad i = 1, \dots, d,$$

which in turn allows us to incorporate a general f.

• Time discrete problem: given the initial condition $\alpha_0 =: \bar{\alpha}_{\tau}(0)$, find $\bar{\alpha}_{\tau} \subset \mathcal{H}^{\tau}_{ad}$ solving

$$ar{m{lpha}}_{ au} = \operatorname*{argmin}_{m{lpha}_{ au} \in \mathcal{H}^{\pi}_{ad}} \mathcal{J}_{ au}(m{lpha}_{ au}), \quad \mathcal{J}_{ au}(m{lpha}_{ au}) = \mathcal{J}^{1}_{ au}(m{lpha}_{ au}) + \mathcal{J}^{2}_{ au}(m{lpha}_{ au}),$$

where

$$\begin{aligned} \mathcal{J}_{\tau}^{1}(\boldsymbol{\alpha}_{\tau}) &+ \mathcal{J}_{\tau}^{2}(\boldsymbol{\alpha}_{\tau}) \\ &= \tau \sum_{n=1}^{N} \frac{1}{2} \sum_{i=1}^{d} \| (\boldsymbol{\alpha}_{\tau}^{n})^{\top} \widehat{\mathbf{P}}_{i}^{n} \boldsymbol{\alpha}_{\tau}^{n} - \widehat{\mathbf{f}}_{i}^{n} \|_{\mathrm{L}^{2}(\widehat{D})}^{2} + \tau \sum_{n=1}^{N} \frac{\lambda}{2\tau^{2}} | \boldsymbol{\alpha}_{\tau}^{n} - \boldsymbol{\alpha}_{\tau}^{n-1} |^{2}. \end{aligned}$$

and

$$\mathcal{H}_{ad}^{\tau} := \{ \boldsymbol{\alpha}_{\tau} \in \mathrm{H}^{1}(0,T) : \boldsymbol{\alpha}_{\tau}|_{[t^{n-1},t^{n}]} \in \mathbb{P}^{1}, \ n = 1, \dots, N \} \cap \mathcal{H}_{ad}.$$

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Convergence of scheme

Theorem. The family of minimizers $\left\{ \bar{\alpha}_{\tau} \right\}_{\tau > 0}$ to the discrete problem is

- uniformly bounded in $H^1(0,T)$.
- it contains a subsequence that converges weakly to $\bar{\alpha}$ in $\mathrm{H}^{1}(0,T)$.

The proof is motivated by Γ convergence.



Strong convergence to local minimizers

Auxiliary problem: For a fixed ε > 0, we construct a family {α^ε_τ}_{τ>0} upon solving the minimization problem

$$\boldsymbol{\alpha}^{\varepsilon}_{\tau} = \operatorname*{argmin}_{\boldsymbol{\alpha}_{\tau} \in \mathcal{H}^{\tau,\varepsilon}_{ad}} \mathcal{J}_{\tau}(\boldsymbol{\alpha}_{\tau}),$$

where $\mathcal{H}_{ad}^{\tau,\varepsilon} = \Big\{ \pmb{\alpha}_{\tau} \in \mathcal{H}_{ad}^{\tau} : \| \Pi_{\tau} \bar{\pmb{\alpha}} - \pmb{\alpha}_{\tau} \|_{\mathrm{L}^{2}(0,T)} \leq \varepsilon \Big\}.$

- Next using the second order sufficient condition we show that {α^ε_τ}_{τ>0} forms a local solution to our discrete problem.
- We conclude by showing that $\| \boldsymbol{\alpha}_{\tau}^{\varepsilon} \bar{\boldsymbol{\alpha}} \|_{\mathrm{H}^{1}(0,T)} \to 0$ as $\tau \to 0$.

This approach is inspired by Casas and Troeltzsch '02.



Problem 2: Discretization

• We consider the following discrete problem: given an initial condition $(\alpha_0, \theta_0) =: (\alpha_{\kappa}(0), \theta_{\kappa}(0))$ find a solution $(\alpha_{\kappa}, \theta_{\kappa}) \in \mathcal{U}_{ad}^{\kappa} \times \mathcal{V}_{ad}^{\kappa}$ to

 $\min_{(\boldsymbol{\alpha}_{\kappa},\boldsymbol{\theta}_{\kappa})\in\mathcal{U}_{ad}^{\kappa}\times\mathcal{V}_{ad}^{\kappa}}\mathcal{F}_{\kappa}(\boldsymbol{\alpha}_{\kappa},\boldsymbol{\theta}_{\kappa})\coloneqq\mathcal{F}_{\kappa}^{1}(\boldsymbol{\alpha}_{\kappa},\boldsymbol{\theta}_{\kappa})+\mathcal{F}_{\kappa}^{2}(\boldsymbol{\theta}_{\kappa})+\mathcal{F}_{\kappa}^{3}(\boldsymbol{\alpha}_{\kappa})+\mathcal{F}_{\kappa}^{4}(\boldsymbol{\theta}_{\kappa})$

where

$$\begin{aligned} \mathcal{F}_{\kappa}^{1}(\boldsymbol{\alpha}_{\kappa},\boldsymbol{\theta}_{\kappa}) &= \sum_{m=1}^{M} \frac{\kappa}{2\theta_{\kappa}^{m}} \sum_{i=1}^{d} \|(\boldsymbol{\alpha}_{\kappa}^{m})^{\top} \widetilde{\mathbf{P}}_{i}^{m} \boldsymbol{\alpha}_{\kappa}^{m} - \boldsymbol{\rho}_{i}'(s^{m}) \boldsymbol{\theta}_{\kappa}^{m} \|_{\mathrm{L}^{2}(\widehat{D})}^{2} \\ \mathcal{F}_{\kappa}^{2}(\boldsymbol{\theta}_{\kappa}) &= \sum_{m=1}^{M} \frac{\beta\kappa}{\theta_{\kappa}^{m}}, \quad \mathcal{F}_{\kappa}^{3}(\boldsymbol{\alpha}_{\kappa}) = \kappa \sum_{m=1}^{M} \frac{\lambda}{2\kappa^{2}} |\boldsymbol{\alpha}_{\kappa}^{m} - \boldsymbol{\alpha}_{\kappa}^{m-1}|^{2}, \\ \mathcal{F}_{\kappa}^{4}(\boldsymbol{\theta}_{\kappa}) &= \kappa \sum_{m=1}^{M} \frac{\eta}{2\kappa^{2}} |\boldsymbol{\theta}_{\kappa}^{m} - \boldsymbol{\theta}_{\kappa}^{m-1}|^{2}. \end{aligned}$$

and admissible sets

$$\mathcal{U}_{ad}^{\kappa} := \left\{ \boldsymbol{\alpha}_{\kappa} \in \mathrm{H}^{1}(0, s_{F}) : \boldsymbol{\alpha}_{\kappa}|_{[s^{m-1}, s^{m}]} \in \mathbb{P}^{1}, m = 1, \dots, M \right\} \cap \mathcal{U}_{ad},$$
$$\mathcal{V}_{ad}^{\kappa} := \left\{ \theta_{\kappa} \in \mathrm{H}^{1}(0, s_{F}) : \theta_{\kappa}|_{[s^{m-1}, s^{m}]} \in \mathbb{P}^{1}, m = 1, \dots, M \right\} \cap \mathcal{V}_{ad}.$$

We again show the weak convergence using Γ-convergence.



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Numerical examples

•
$$\Omega = B_1(0,0) \subset \mathbb{R}^2$$
, $\widehat{D} = B_{0.2}(-0.75,0)$
• $\boldsymbol{x}_{k+1} = 1.2(\cos(k\pi/4), \sin(k\pi/4))$
• $\widehat{\mathbf{d}}_{k+1} = (\cos(k\pi/4), \sin(k\pi/4)), k = 0, \dots, 7 \ (n_p = 8)$





Problem 1: approximate $\mathbf{f}_1(\boldsymbol{x},t) = (1,0)^{\top}$







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Magnetic force at t = 0.0125, 0.5, and 1.



Dipoles on the left (dipoles 4, 5 and 6) have small intensities at initial times. This is expected because $D_{1,t}$ is close to the boundary of Ω , where **H** is large, thus it is difficult for dipoles 4, 5 and 6 to "push" in the **f**₁ direction.

Problem 1: approximate $\mathbf{f}_1(\boldsymbol{x},t) = (1,0)^\top$

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Problem 1: approximate $\mathbf{f}_2(\boldsymbol{x},t) = (\sin(\pi(1-t)), -\cos(\pi(1-t))^{\top})$





Problem 1: approximate $\mathbf{f}_2(\boldsymbol{x},t) = (\sin(\pi(1-t)), -\cos(\pi(1-t))^{\top})$



Magnetic force at t = 0.015, 0.375 and 0.75.





Problem 1: approximate $\mathbf{f}_2(\boldsymbol{x},t) = (\sin(\pi(1-t)), -\cos(\pi(1-t))^{\top})$

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Problem 2

• Let the curve C is parameterized by

$$oldsymbol{
ho}(s) = oldsymbol{x}_I + s rac{(oldsymbol{x}_F - oldsymbol{x}_I)}{\|oldsymbol{x}_F - oldsymbol{x}_I\|}, \qquad s \in [0, 0.75]$$

with
$$\boldsymbol{x}_I = (0, -0.75)$$
 and $\boldsymbol{x}_F = (0, 0)$.
 $(\boldsymbol{\alpha}_*, \boldsymbol{\alpha}^*, \theta_*, \theta^*) = (-1, 1, 10^{-10}, 10)$.
 $\boldsymbol{\beta} = 10^{-1}, \, \boldsymbol{\lambda} = 10^{-6}, \text{ and } \boldsymbol{\eta} = 10^{-4})$.



Application: concentration transport

- ▶ **MNPs**. We assume a concentration of magnetic nanoparticles confined in a domain $\Omega \subset \mathbb{R}^d, d = 2, 3$.
- Drug concentration is evolved using

$$\begin{aligned} \frac{\partial c}{\partial t} + \operatorname{div} \left(-A\nabla c + c\mathbf{u} + \gamma_1 cf(\mathbf{H}) \right) &= 0 \quad \text{in } \Omega \times (0, T) \\ c &= 0 \quad \text{on } \partial\Omega \times (0, T) \qquad c(x, 0) = c_0 \quad \text{in } \Omega \\ \operatorname{curl} \mathbf{H} &= \mathbf{0} \quad \text{in } \Omega \qquad \operatorname{div} (\mathbf{H}) = 0 \quad \text{in } \Omega \end{aligned}$$

where $A=10^{-3}$ is a diffusion coefficient matrix, ${\bf u}$ is a fixed velocity vector and f is the Kelvin force.

Goal. Move c_0 from one subdomain to another (desired location) using the magnetic force f while minimizing the spreading.



Application: concentration transport

We solve the parabolic problem with magnetic force given by Problem 1 with f_1 .

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Conclusions

We have approximated a vector field by using the Kelvin force. In particular, we study two problems:

- Fixed final time
- Unknown final time
- We prove the existence of solution and using second order sufficient conditions we show the local uniqueness.
- Motivated by Γ-convergence we show the H¹-weak convergence of the time-discrete problems.
- In presence of second order sufficient condition, a H¹-strong local convergence result is proved for Problem 1.
- As an application, we study the control of magnetic nanoparticles as those used in magnetic drug delivery. The optimized Kelvin force is used to transport the drug to a desired location.

